

# Probing Quark-Gluon Plasma with Bottom Quark Jets at sPHENIX - Theory

Expertise from T-2, T-5, CCS-7

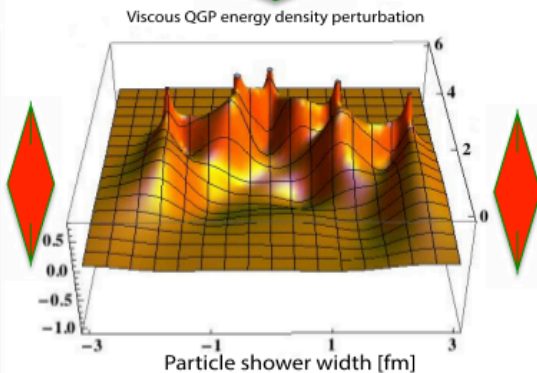
Daligault, Gupta, Kang, Lee, Vitev (co-PI), Yoon

Theory budget ~ \$600K/year



*Perturbative QCD/SCET and jet simulations:* most precise b-jet theory in proton collisions, new theory for heavy ion collisions, b-jet substructure, b-jet tomography of the QGP

*QMD simulations:*  
transport properties of plasmas, stopping power for heavy particles



*Lattice QCD:*  
EoS, input for hydro, charge number fluctuations near the phase transition

*Experiment:* Tracker design, prototype construction, jet finder development, ongoing and improved PHENIX and STAR BES II analyses

## b-jet energy loss

*pQCD / SCET*

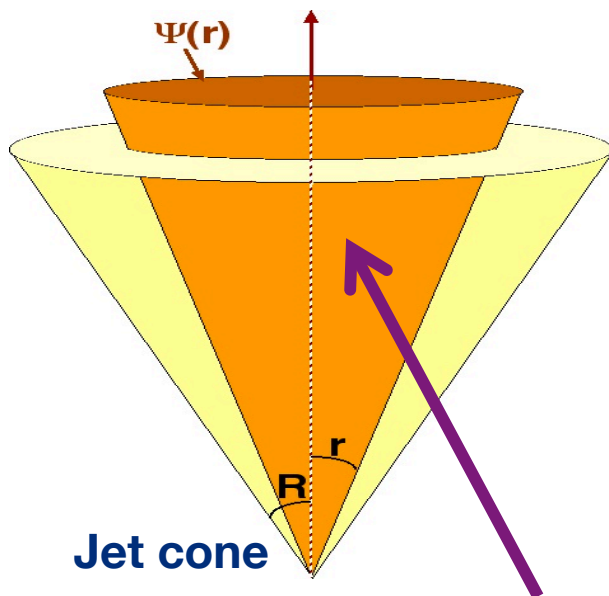
*MD simulations /  $dE/dx$*

*Lattice QCD / Hydro*

# Precision b-jets X sections and substructure<sup>2</sup>

Kang, Lee, Vitev

- Powerful soft-collinear effective field theory (SCET) methods. Vast improvements in accuracy of jet cross sections in  $e^+e^-$ ,  $ep$ , and  $pp$  collisions

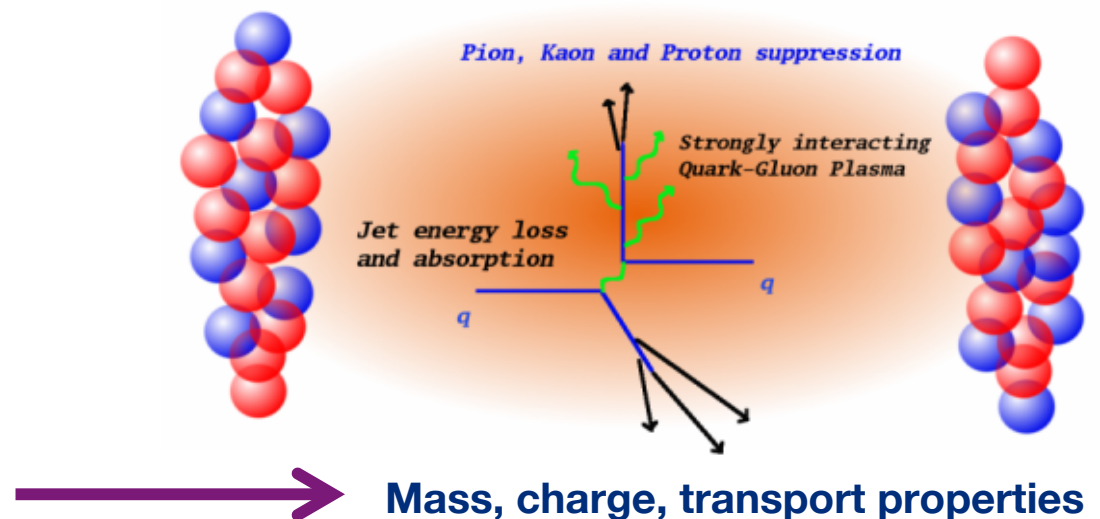


**Evaluate jet substructure observables (shapes, fragmentation functions) and their modification in the QGP**

**Improved formulations of SCET, compute b-jet cross sections with the highest accuracy to date**

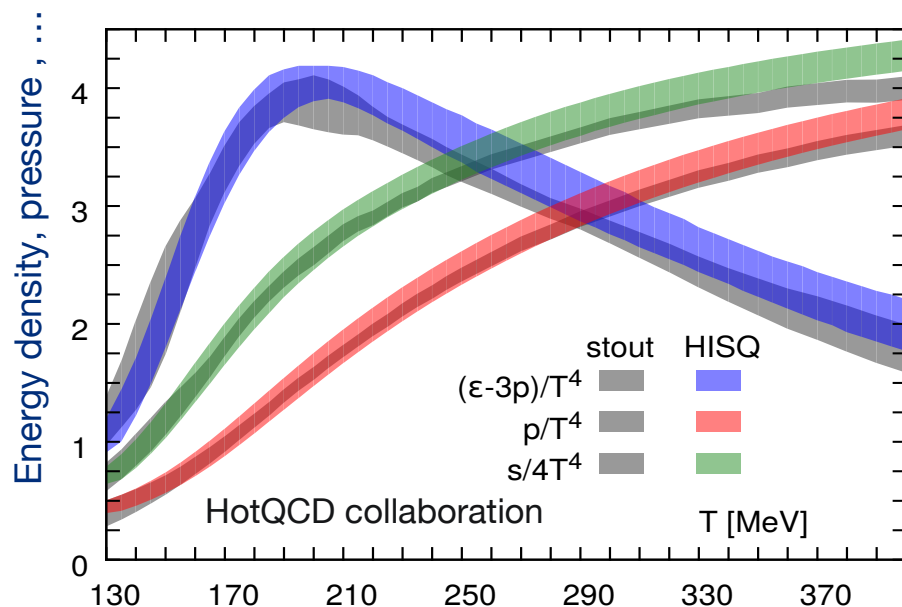
Phys. Rev. D90 (2014) 094503

**Develop first-principles theory of heavy quark propagation in nuclear matter and the process of shower formation**



# Lattice QCD EoS and Charge Fluctuations

3



Phys. Rev. D90 (2014) 094503

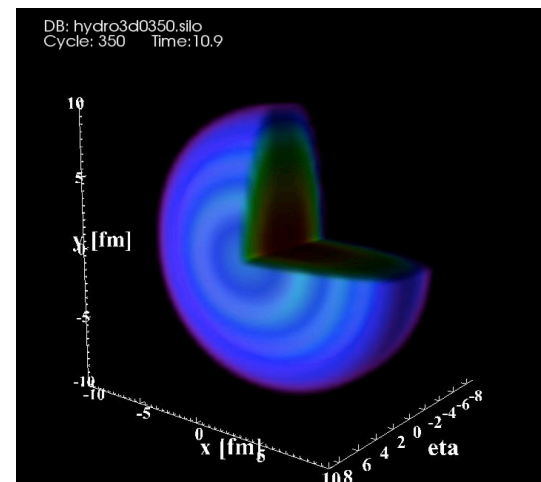
- Phase transitions near the critical point characterized by fluctuations (charge number fluctuation for QGP to ordinary nuclear matter)  $\sim (\pi^+ - \pi^-)$

**Evaluate charge fluctuations on the lattice and compare our results to measurements from Beam Energy Scan II at RHIC**

Gupta, Yoon

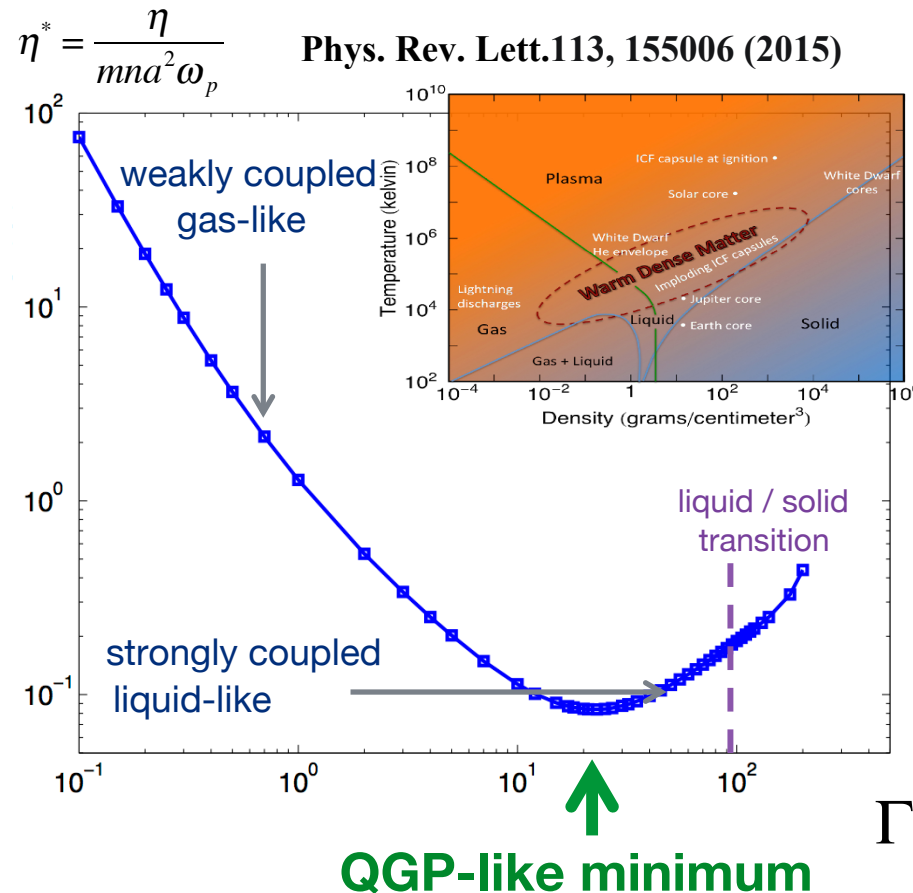
- Incorporate the current state-of-the-art EoS in hydrodynamic simulations, describe time, temperature evolution of the QGP

**Accurate medium dynamics for b-jet quenching simulations will be available**



Hydro simulation of the QGP energy density at  $10^{-23}$  s

# dE/dx in Strongly-Coupled Plasmas



Daligault, Vitev

- The strongly coupled nature of the QGP makes it tantalizingly similar to warm dense matter (WDM). Microscopic MD simulations

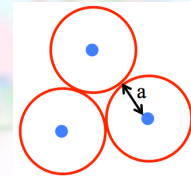
Coupling strength:  $\Gamma = \frac{PE}{KE}$

$\Gamma > 1$

Interactions among particles dramatically affect their dynamics

In WDM,

$$\Gamma = \frac{Q^2 / a}{k_B T}$$



In QGP,

$\Gamma = \text{few}$

**Perform molecular dynamics (MD) simulations of stopping power of charged particles in WDM near the viscosity minimum**

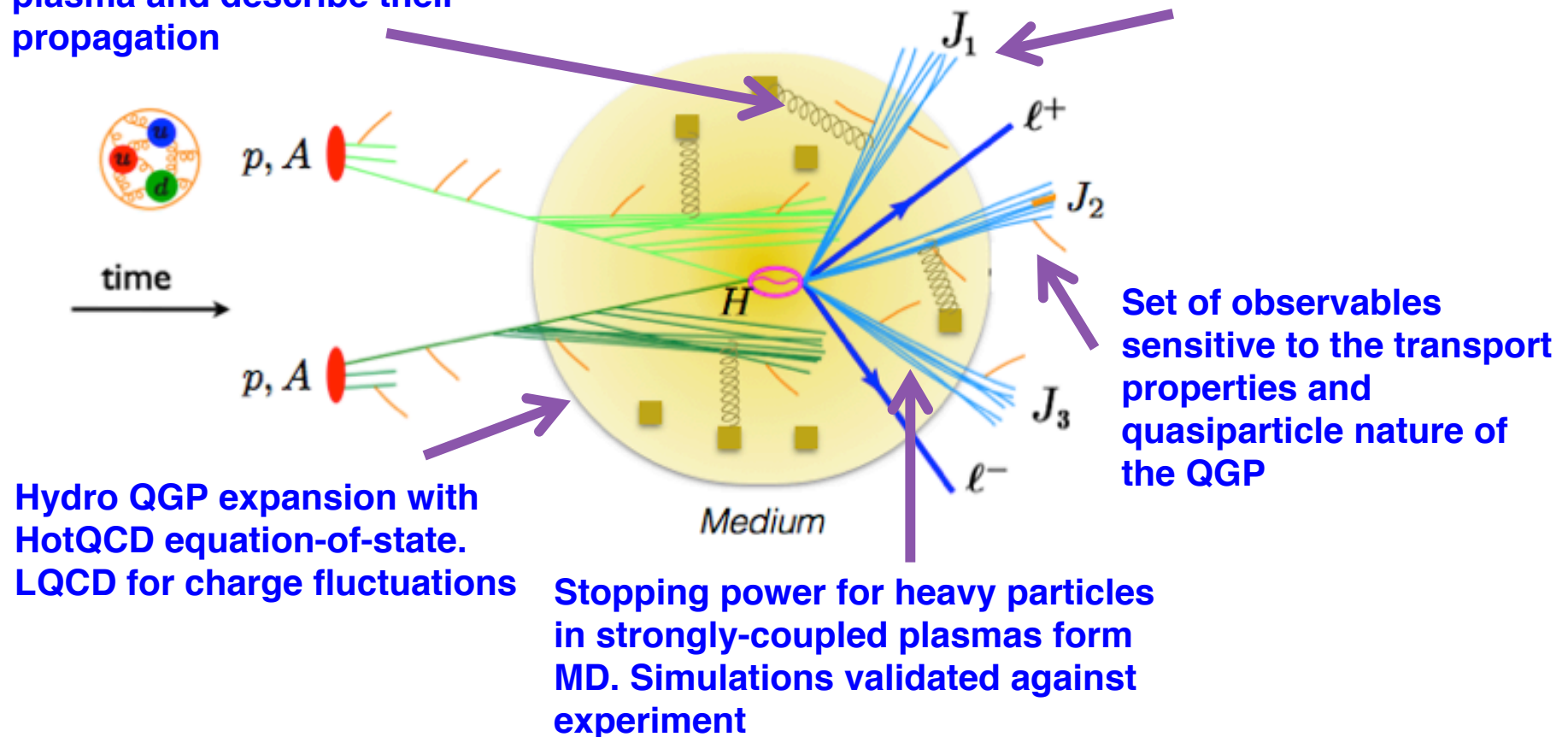
**Obtain much-needed physical insights and theoretical guidance for b-jet stopping power phenomenology in the QGP**

# Theory Deliverables: a Unified Picture

5

First-principles effective theory to couple b-jets to the plasma and describe their propagation

Most accurate resummed calculations of b-jet cross sections and b-jet substructure in proton-proton collisions at RHIC and LHC

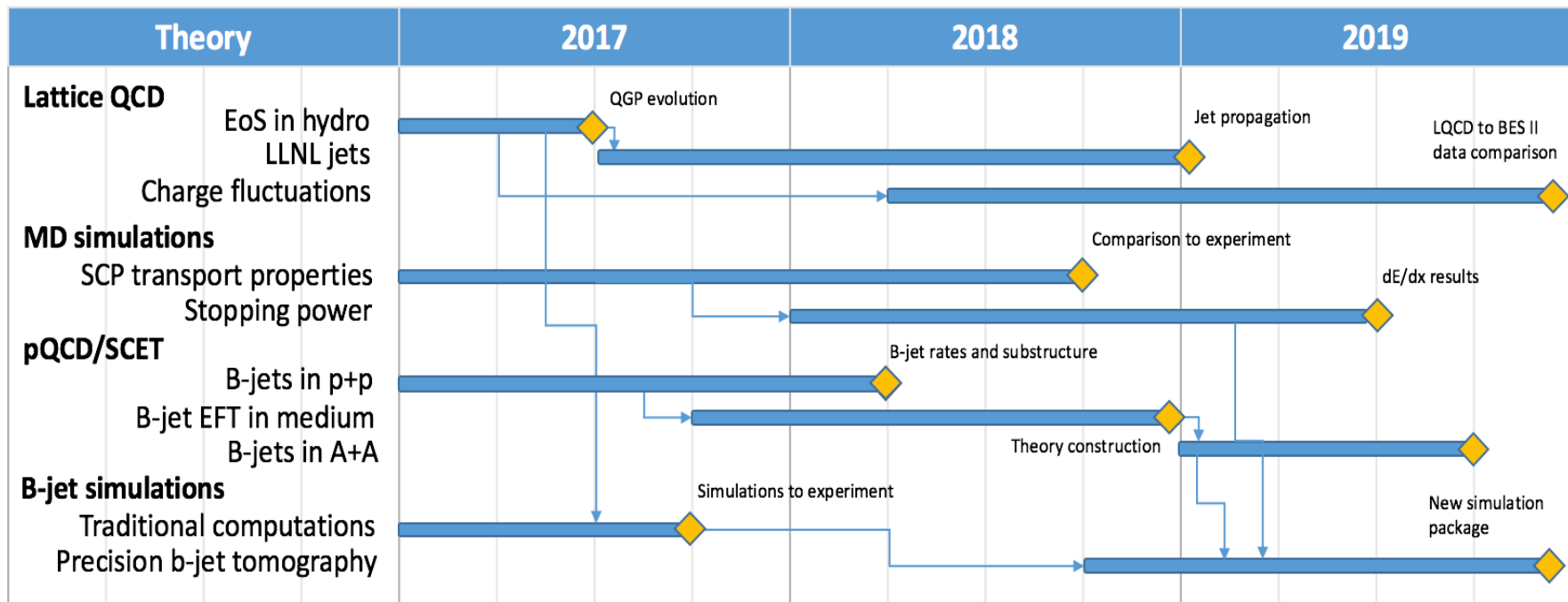


Package for precision b-jet tomography at sPHENIX and predictions to experiment



# Theory Timetable

6



- Individual theory discussion on weekly basis, keeping track of progress
- Well-defined responsibilities, researchers matched to tasks
- Monthly combined theory and experiment meetings
- Progress on items with milestones this year (hydro simulations with LQCD EoS and b-jet simulations with energy loss)
- One item completed ahead of schedule – EFT for heavy quark propagation in the QGP

# One Theory Highlight

## Develop an effective theory of b-jet propagation in matter

FY 2018

### Earlier developments

- Develop first-principles theory of heavy quark propagation in nuclear matter and the process of shower formation

Phys. Lett. B564 (2003) 231-234

JHEP 1106 (2011) 080

It was possible to carry this simplification since we realized the sectors of the theory decouple, very explicit in the hybrid gauge

## LO SCET Lagrangian

$$\mathcal{L}_0 = \sum_{\tilde{p}, \tilde{p}', \tilde{q}} e^{-ix \cdot \mathcal{P}} \bar{\xi}_{n,p'} \left[ i n \cdot D + (\not{\mathcal{P}}_{\perp} + g A_{n,q}^{\perp}) W_n \frac{1}{\not{\mathcal{P}}} W_n^{\dagger} (\not{\mathcal{P}}_{\perp} + g A_{n,q'}^{\perp}) \right] \frac{\not{n}}{2} \xi_{n,p} + \mathcal{L}_m$$

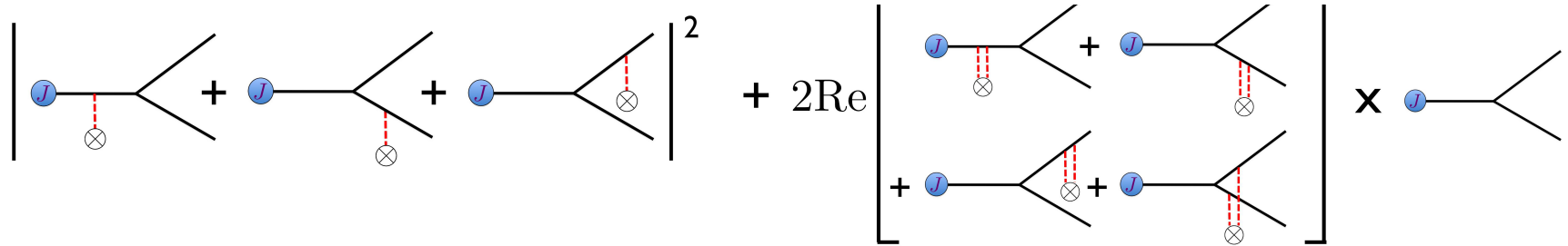
## The mass term

$$\mathcal{L}_m = \sum_{\tilde{p}, \tilde{p}', \tilde{q}} e^{-ix \cdot \mathcal{P}} \left[ m \bar{\xi}_{n,p'} \left[ (\not{\mathcal{P}}_{\perp} + g A_{n,q}^{\perp}), W_n \frac{1}{\not{\mathcal{P}}} W_n^{\dagger} \right] \frac{\not{n}}{2} \xi_{n,p} - m^2 \bar{\xi}_{n,p'} W_n \frac{1}{\not{\mathcal{P}}} W_n^{\dagger} \frac{\not{n}}{2} \xi_{n,p} \right]$$

The mass  $m/p^+ \sim \lambda$  in  $\text{SCET}_M$  already counts as the small parameter. So any term that will include mass and the Glauber field will be power suppressed

# Necessary Ingredient - Massive Splittings

## Diagrams corresponding to first order in opacity

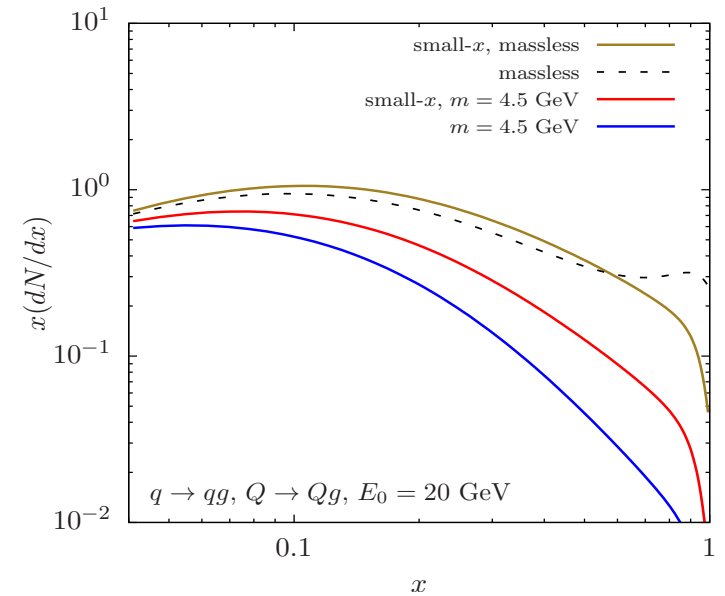


- First recover the massive vacuum splittings
- Evaluate the 3 massive in medium splittings

We have input for description of the in-medium parton showers of heavy quarks

One example. Can be evaluated numerically

$$\begin{aligned}
 \left( \frac{dN^{\text{med}}}{dx d^2k_{\perp}} \right)_{Q \rightarrow Qg} &= \frac{\alpha_s}{2\pi^2} C_F \int \frac{d\Delta z}{\lambda_g(z)} \int d^2q_{\perp} \frac{1}{\sigma_{el}} \frac{d\sigma_{el}^{\text{med}}}{d^2q_{\perp}} \left\{ \left( \frac{1 + (1-x)^2}{x} \right) \left[ \frac{B_{\perp}}{B_{\perp}^2 + \nu^2} \right. \right. \\
 &\times \left( \frac{B_{\perp}}{B_{\perp}^2 + \nu^2} - \frac{C_{\perp}}{C_{\perp}^2 + \nu^2} \right) (1 - \cos[(\Omega_1 - \Omega_2)\Delta z]) + \frac{C_{\perp}}{C_{\perp}^2 + \nu^2} \cdot \left( 2 \frac{C_{\perp}}{C_{\perp}^2 + \nu^2} - \frac{A_{\perp}}{A_{\perp}^2 + \nu^2} \right. \\
 &- \left. \left. \frac{B_{\perp}}{B_{\perp}^2 + \nu^2} \right) (1 - \cos[(\Omega_1 - \Omega_3)\Delta z]) + \frac{B_{\perp}}{B_{\perp}^2 + \nu^2} \cdot \frac{C_{\perp}}{C_{\perp}^2 + \nu^2} (1 - \cos[(\Omega_2 - \Omega_3)\Delta z]) \right. \\
 &+ \frac{A_{\perp}}{A_{\perp}^2 + \nu^2} \cdot \left( \frac{D_{\perp}}{D_{\perp}^2 + \nu^2} - \frac{A_{\perp}}{A_{\perp}^2 + \nu^2} \right) (1 - \cos[\Omega_4 \Delta z]) - \frac{A_{\perp}}{A_{\perp}^2 + \nu^2} \cdot \frac{D_{\perp}}{D_{\perp}^2 + \nu^2} (1 - \cos[\Omega_5 \Delta z]) \\
 &+ \left. \left. \frac{1}{N_c^2} \frac{B_{\perp}}{B_{\perp}^2 + \nu^2} \cdot \left( \frac{A_{\perp}}{A_{\perp}^2 + \nu^2} - \frac{B_{\perp}}{B_{\perp}^2 + \nu^2} \right) (1 - \cos[(\Omega_1 - \Omega_2)\Delta z]) \right] \right\} \\
 &+ x^3 m^2 \left[ \frac{1}{B_{\perp}^2 + \nu^2} \cdot \left( \frac{1}{B_{\perp}^2 + \nu^2} - \frac{1}{C_{\perp}^2 + \nu^2} \right) (1 - \cos[(\Omega_1 - \Omega_2)\Delta z]) + \dots \right] \Bigg\}
 \end{aligned}$$



Largest mass effects in the sPHENIX b-jet acceptance range

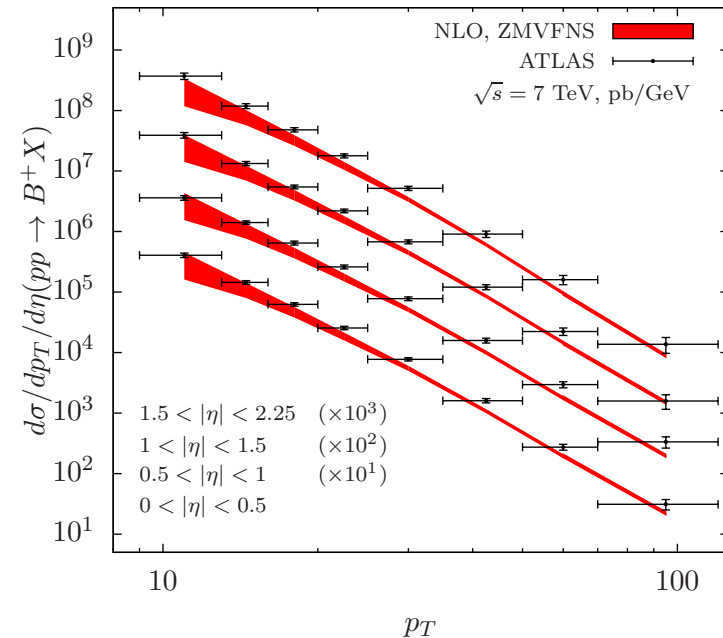
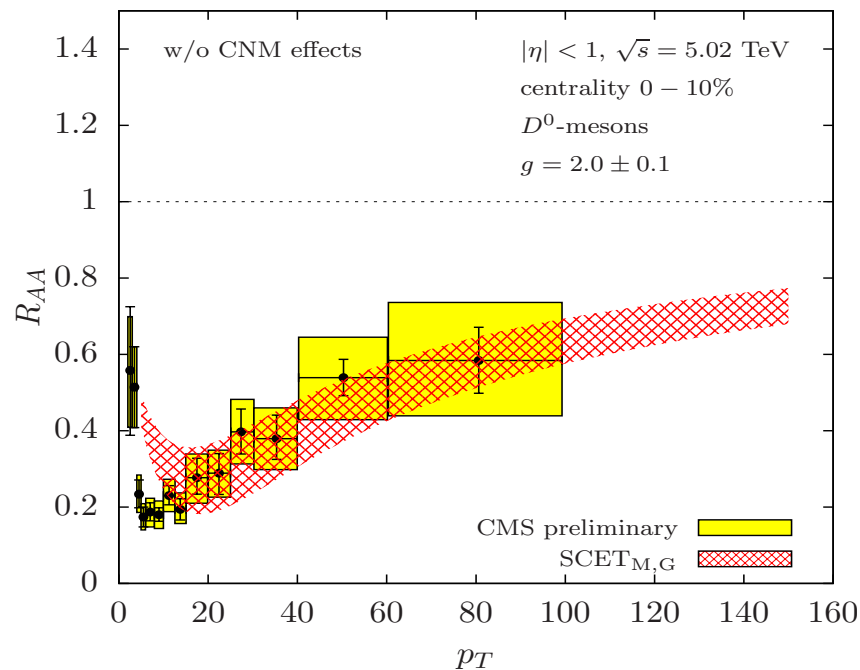


# Application Consistent with NLO

Two advances – a) including gluon fragmentation contribution to heavy flavor, b) going beyond energy loss

$$\frac{d\sigma_{pp}^H}{dp_T d\eta} = \frac{2p_T}{s} \sum_{a,b,c} \int_{x_a^{\min}}^1 \frac{dx_a}{x_a} f_a(x_a, \mu) \int_{x_b^{\min}}^1 \frac{dx_b}{x_b} f_b(x_b, \mu) \times \int_{z_c^{\min}}^1 \frac{dz_c}{z_c^2} \frac{d\hat{\sigma}_{ab}^c(\hat{s}, \hat{p}_T, \hat{\eta}, \mu)}{dvdz} D_c^H(z_c, \mu),$$

- Gluon fragmentation plays an important role ~50%



- First cascade contribution

$$d\sigma_{PbPb}^H = d\sigma_{pp}^{H,NLO} + d\sigma_{PbPb}^{H,med}$$

This leaves us time for a stretch goal – include collisional energy losses in the SCET framework. This can go well through the second year

# Conclusions

10

- Theory is an integral part of this LDRD DR
- There are project management mechanisms in place and ways to track theory progress.  
Milestones, clear responsibilities assigned
- At present, the theory part is on schedule  
[Hydro code installed, tests being run]  
[HF parton shower in the soft emission limit simulated]
- One milestone completed ahead of schedule  
[Effective theory for open heavy propagation in matter]